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## Hydroxyapatite substituted by transition metals: experiment and theory

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Bioceramics are versatile materials for hard tissue engineering. Hydroxyapatite (HA) is a widely studied biomaterial for bone grafting and tissue engineering applications. The crystal structure of HA allows for a wide range of substitutions, which allows for tailoring materials properties. Transition metals and lanthanides are of interest since substitution in HA can result in magnetic properties. In this study, experimental results were compared to theoretical calculations of HA substituted with a transition metal. Calculation of a 10 atomic percent substitution of a transition metal ion  $Mn^{2+}$ ,  $Fe^{2+}$ , and  $Co^{2+}$  substituted HA samples lead to magnetic moments of 5, 4, and 3 Bohr magnetons, respectively. Hydroxyapatite substituted by transition metals (MHA) was fabricated through an ion exchange procedure and characterized with X-ray diffraction, Fourier transform infra-red spectroscopy (FTIR), X-ray photoelectron spectroscopy, and vibrating sample magnetometer, and results were compared to theoretical calculations. All the substitutions resulted in phase-pure  $M^{2+}HA$  with lattice parameters and FTIR spectra in good agreement with calculations. Magnetic measurements revealed that the substitution of  $Mn^{2+}$  has the greatest effect on the magnetic properties of HA followed by the substitution of  $Fe^{2+}$  and then  $Co^{2+}$ . The present work underlines the power of synergistic theoretical–experimental work in guiding the rational design of materials.

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### 1. Introduction

Hydroxyapatite (HA) is a commonly used bone grafting material and is a well-documented biomaterial known for its osteoconductivity.<sup>1,2</sup> The main mineral phase of bone is HA,  $Ca_{10}(PO_4)_6(OH)_2$ , and its bioactivity can be modified through substitutions of other anions, cations, or functional groups. Well studied and commonly found trace elements in bone are  $Mg^{2+}$ ,  $Si^{2+}$ ,  $CO_3^{2-}$ , and  $SiO_4^{4-}$ ; these elements have been shown to render HA more bioactive compared to stoichiometric HA.<sup>3–6</sup> The hexagonal crystal structure of HA allows for a wide variety of substitutions that can alter bioactivity but also physical properties as well.<sup>7</sup> Other possible substituents which are less studied are transition metals of the 3d block, which are necessary for normal biological function.<sup>8,9</sup>

One unique characteristic of some transition metals is their magnetic response, one that can be modified and tuned in HA, which is diamagnetic, *via* substitutions. Thus, through substituting with magnetic ions HA with magnetic properties can

be achieved resulting in a biocompatible and magnetic biomaterial. There is a wide range of biomedical applications for biocompatible and magnetic materials such as drug delivery, cell separation, magnetic resonance imaging, and hyperthermia applications.<sup>10–13</sup> Iron oxide nanoparticles are currently used for these applications but there have been concerns over acute toxicity.<sup>14</sup> Since HA is biocompatible and biodegradable, application of magnetic HA will mitigate these concerns.

The most commonly used metal for imparting HA with magnetic properties is iron.<sup>15–18</sup> The choice of iron as a dopant stems from its high daily dietary need and its large magnetic moment. Other transition metals possess the potential and have rendered HA magnetic, such as cobalt,<sup>19</sup> but concerns over toxicity has limited the research despite the fact that cobalt is required in trace amounts for vitamin B12 function.<sup>20</sup> Another suitable transition metal that has received little attention is manganese, which has been shown to influence bone remodeling and is used as a supplement to retard bone loss.<sup>21,22</sup> However, initial studies have reported that manganese substitution does not affect the magnetic properties of HA.<sup>23</sup> Many elements have the potential to yield magnetic HA with varying levels of bioactivity and magnetization.

The time to synthesize and characterize all possible magnetic HAs is laborious and time consuming. The use of density functional theory (DFT) calculations to identify promising new

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materials helps to guide and expedite experimental work. Computational efforts coupled with experimental work has provided information on site preference, structural modifications, and possible mechanisms of substitution in Pb, Cd, Sr, Mg, Zn, Ti, Si,  $\text{CO}_3^{2-}$ , and F.<sup>24–34</sup> In Si substituted HA, DFT has been used to study defect chemistry with results suggesting two different charge compensation mechanisms (1) the formation of  $\text{HSiO}_4$  groups under hydrating conditions and (2) under dehydration the formation of OH vacancies, again DFT results coincided well with experimental data.<sup>31,32</sup> Material property changes have been successfully predicted with DFT in calcium deficient hydroxyapatite where Bhat *et al.* calculated a decrease of 30% in elastic properties due to calcium vacancies which coincided with observed experimental results.<sup>35</sup> Optical properties have also been calculated for the  $\text{Ca}_{10}(\text{PO}_4)_6\text{X}_2$  system where X is OH, F, Cl, or Br with DFT and only minor differences in optical properties were calculated.<sup>36</sup> Theoretical work on magnetic HA systems have focused on the calculation of magnetic spectroscopic parameters for electron paramagnetic resonance (EPR) and Mössbauer spectras<sup>37–40</sup> Jiang *et al.* used DFT to analyze electronic structure and compute magnetic spectroscopic parameters for EPR and Mössbauer spectras of  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  substituted HA.<sup>37</sup> The comparison of experimental spectrums with computed spectroscopic parameters allowed for the assignment of site occupancy of  $\text{Fe}^{2+}$  and  $\text{Fe}^{3+}$  in the HA crystal lattice.<sup>37</sup> While EPR studies combined with DFT calculations have also helped understand interactions between nitrate and  $\text{Mn}^{2+}$  ions and nitrate and carbonate ions in the HA lattice where the nitrate ions are the paramagnetic center of interest.<sup>38–40</sup>

Little theoretical work has been done on the magnetic properties of HA substituted by divalent transition metals (MHA). Using DFT, we determined the most promising transition metals that would yield the highest magnetic moment, followed by synthesis of the substituted HAs. Herein we investigate the substitution of divalent transition metal ions:  $\text{Mn}^{2+}$ ,  $\text{Fe}^{2+}$ , and  $\text{Co}^{2+}$  in HA from both theoretical and experimental approaches.

## 2. Materials and methods

### 2.1 Materials

Ammonium phosphate dibasic (99+%), iron chloride tetrahydrate (99+%), cobalt chloride hexahydrate (analysis grade), sulfuric acid (95–98%), and potassium bromide (IR grade) were purchased from Acros Organics. Calcium nitrate (99+%), manganese chloride tetrahydrate (99+%), lanthanum chloride heptahydrate (99+%), nitric acid (69.2%), and ammonium hydroxide (29.45%) were purchased from Fisher Scientific.

### 2.2 HA and MHA synthesis

HA was synthesized *via* a wet chemical method. An aqueous solution of ammonium phosphate dibasic (80 mM) was added drop-wise to an aqueous calcium nitrate solution (225 mM) at room temperature with the calcium to phosphate ratio maintained at 1.667. Ammonium hydroxide was used to raise the pH of both solutions above 11 prior to mixing the two solutions.

Once the two solutions were completely mixed, the temperature was raised to 95 °C and maintained at the temperature for five hours. Then the solution was cooled to room temperature and aged for two days under constant stirring. The precipitates were collected *via* centrifugation and washed 3 times with de-ionized, 2 times with ethanol and vacuum dried.

MHA samples were fabricated through an ion-exchange procedure. A salt solution of the desired transition metal ( $\text{MnCl}_2$ ,  $\text{FeCl}_2$ , or  $\text{CoCl}_2$ ) was prepared at a concentration of 0.02 M, and the pH of the solution was adjusted to 2.7 with dilute sulfuric acid. Previously synthesized HA was immersed in the ion-exchange solution at 0.5 g  $\text{dL}^{-1}$  and sonicated for one hour at room temperature with intermittent stirring. The ion-exchanged HA was collected by centrifugation, washed 3 times with DIW, and vacuum dried.

### 2.3 Characterization

Chemical analysis of the calcium and transition metal was determined by atomic absorption spectroscopy (Perkin Elmer Atomic Adsorption Spectrometer 3100). In brief, 15 mg of MHA powder was dissolved in 200  $\mu\text{L}$  of concentrated nitric acid (15.9 M). A 2 M nitric acid solution with 5% (w/v) lanthanum chloride was then used to raise the total volume to 1 mL. The total volume was diluted to 20 mL with 0.1 M nitric acid and dilutions of the solution were used for chemical analysis.

Powder X-ray diffraction (XRD) patterns were acquired with a Bruker D2 phaser X-ray diffractometer operating at 30 kV and 10 mA using Cu K $\alpha$  radiation. Diffraction patterns were collected over  $2\theta$  between 10–120° with a step size of 0.01° and a scan rate of 4 s per step. Powder patterns were refined using the Rietveld method with the general structural analysis system (GSAS) and the EXPGUI interface.<sup>41,42</sup>

FTIR spectra were acquired on a Nicolet Magna 560 FTIR spectrometer. Specimens were prepared using the potassium bromide method. In brief, powder samples were diluted to one part per million in KBr, grinded into a fine powder and then pressed into a transparent pellet. The FTIR spectra were acquired over the range of 4000–400  $\text{cm}^{-1}$  with a resolution of 4  $\text{cm}^{-1}$  and a total of 32 scans.

In order to assess the surface composition and oxidation state of the transition metal dopant, X-ray photoelectron spectroscopy (XPS) was employed using a VG X-ray Photoelectron Spectrometer with an Al K $\alpha$  radiation source. High resolution XPS spectrums were acquired of the transition metal 2p and O 1s peaks using a pass energy of 50 eV and a step size of 0.1 eV. Spectrums were calibrated by setting the adventitious carbon peak to 284.7 eV, and CasaXPS was used to deconvolute the spectra.

Magnetic measurements were performed using a vibrating sample magnetometer (VSM) attached to the evercool physical property measurement system (from Quantum Design). Magnetization measurements were carried out at room temperature and with applied DC magnetic fields of up to 1 Tesla.

### 2.4 Theoretical calculations

Our DFT calculations were performed with the Vienna *ab initio* simulation package (VASP),<sup>43</sup> using the Perdew–Burke–Ernzerhof

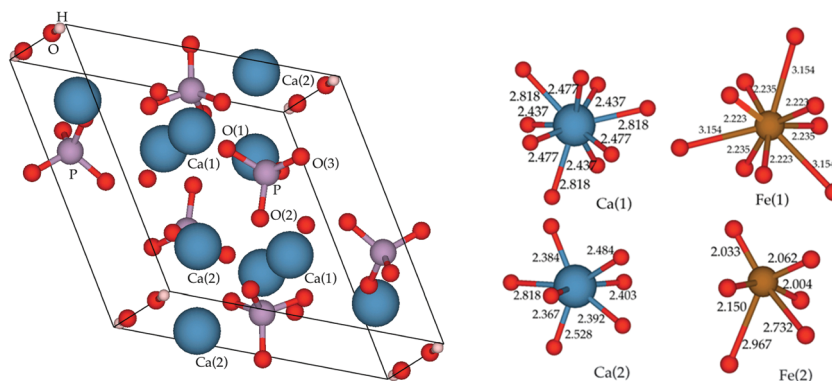


Fig. 1 On the left is the unit cell of hydroxyapatite from the calculated atomic positions. To the right of the unit cell is the local environment of Ca in pure HA and Fe in FeHA. Bond distances are in Å.

(PBE)<sup>44</sup> generalized gradient approximation (GGA) and a plane wave energy cutoff of 750 eV. Dense Monkhorst–Pack  $k$ -point meshes<sup>45</sup> of  $2 \times 2 \times 2$  were used for sampling the Brillouin zone of HA. All atoms were allowed to relax to their equilibrium positions until atomic forces are smaller than  $0.01 \text{ eV } \text{\AA}^{-1}$ .

### 3. Hydroxyapatite crystal structure

HA occurs in a hexagonal crystal structure (space group  $P6_3/m$ ) with 44 atoms per unit cell and lattice parameters  $a = b \neq c$  and  $\alpha = \beta = 90^\circ$ ,  $\gamma = 120^\circ$ . The unit cell and the local environment of Ca in pure HA is depicted in Fig. 1. There are two distinct Ca sites where cation exchange may occur. The Ca(1) site is surrounded by nine oxygen ions from the surrounding phosphate tetrahedron and forms a distinct Ca channel parallel to the  $c$ -axis. The Ca(2) site is surrounded by six oxygen ions from five phosphate groups and the hydroxyl ion. The Ca(2) ions form triangular bases staggered by  $60^\circ$  above and below the OH ion. The OH ion similar to Ca(1) forms a distinct channel parallel to the  $c$ -axis. Based on the distance and the coordination number of oxygen and calcium, three oxygen atoms are defined as O(1), O(2) and O(3). The oxygen in  $\text{OH}^-$  is defined as O<sub>H</sub>. Thus the formula for HA can also be written as  $\text{Ca}_{10}\text{Ca}_{11}\text{O}_{20}(\text{PO}_4)_6(\text{OH})_2$ .

The theoretical lattice constants from DFT calculation are  $a = b = 9.55 \text{ \AA}$  and  $c = 6.94 \text{ \AA}$ , which are in good agreement with the previously reported experimental values of HA:  $a = b = 9.43 \text{ \AA}$ , and  $c = 6.88 \text{ \AA}$ .<sup>46</sup> We also compare atomic positions and relaxed bond lengths in the unit cell with experimental and past DFT results for undoped HA in Tables 1 and 2. Our results are very close to the experimental data.<sup>47</sup> Based on Fig. 1, we note that there are six nearest neighboring oxygen atoms and three second nearest adjacent oxygen atoms for Ca(1), while Ca(2) has six nearest neighboring oxygen atoms and only one second nearest adjacent oxygen atom.

### 4. Results and discussion

The phase purity of the as-synthesized HA and MHA was evaluated using X-ray diffraction. Powder diffraction patterns of these

Table 1 Calculated atomic fractional coordinates of HA compare to literature values

Atom	This work			Other DFT <sup>a</sup>			Exp <sup>b</sup>		
	$x$	$y$	$z$	$x$	$y$	$z$	$x$	$y$	$z$
Ca(1)	0.333	0.667	0.004	0.333	0.667	0.003	0.333	0.667	0.001
Ca(2)	0.251	0.999	0.256	0.251	1	0.25	0.246	0.993	0.25
P	0.399	0.369	0.248	0.4	0.369	0.252	0.400	0.369	0.25
O(1)	0.332	0.486	0.251	0.333	0.487	0.253	0.329	0.484	0.25
O(2)	0.588	0.464	0.245	0.588	0.463	0.247	0.589	0.466	0.25
O(3)	0.34	0.257	0.068	0.338	0.256	0.073	0.348	0.259	0.073
OH	0	0	0.246	0	0	0.285	0	0	0.25

<sup>a</sup> Ref. 55. <sup>b</sup> Ref. 46.

Table 2 Calculated bond lengths (Å) and bond angles ( $^\circ$ ) for HA compared to literature values

	This work	Other DFT <sup>a</sup>	Exp <sup>b</sup>
Bond length			
P–O(1)	1.550	1.541	1.533
P–O(2)	1.563	1.554	1.544
P–O(3)	1.551	1.545	1.514
Ca(1)–O(1)	2.437	2.425	2.416
Ca(1)–O(2)	2.477	2.469	2.449
Ca(1)–O(3)	2.817	2.827	2.802
Ca(2)–O(1)	2.818	2.798	2.712
Ca(2)–O(2)	2.366	2.357	2.356
Ca(2)–O(3)	2.391	2.367	2.367
Bond angles			
O(1)–P–O(2)	111		
O(1)–P–O(3)	111.8		
O(3)–P–O(2)	107.9		
O(1)–Ca(2)–O(2)	47.3		
O(1)–Ca(1)–O(2)	28.5		

<sup>a</sup> Ref. 55. <sup>b</sup> Ref. 46.

samples are depicted in Fig. 2 along with the computed XRD patterns from the theoretical model, Fig. 3. The as-synthesized HA is identified with JCPDF 09-432 and the substituted HA do not exhibit any extraneous peaks, suggesting a single phase in all the powders. Rietveld refinements were performed on the synthesized powders using the refined atomic positions reported by Veselinovic *et al.*<sup>48</sup> The refined lattice parameters of HA and MHA are summarized in Table 3, which are in good

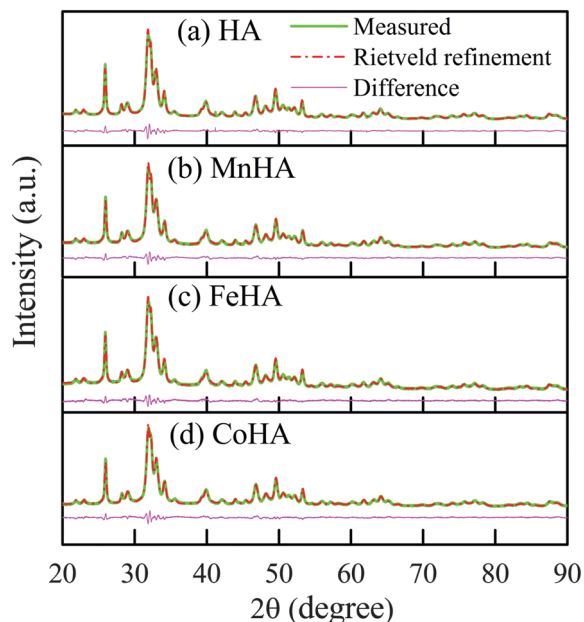


Fig. 2 Powder XRD patterns of the synthesized HA and MHA: HA (a), MnHA (b), FeHA (c), and CoHA (d).

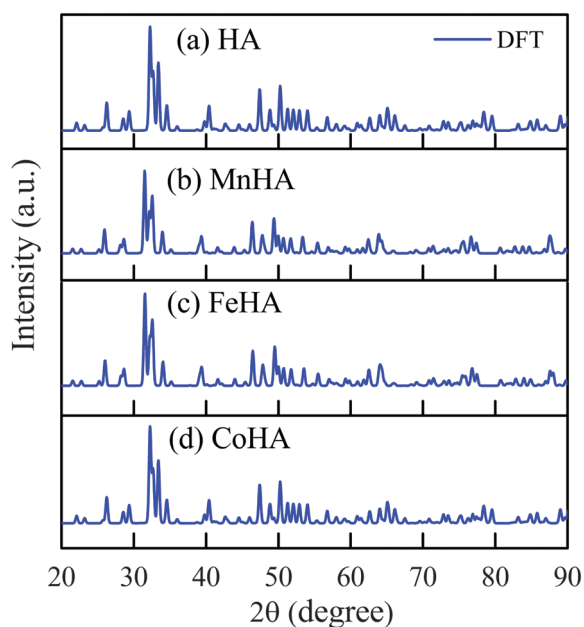


Fig. 3 Calculated XRD patterns from DFT calculations of HA and MHA: HA (a), MnHA (b), FeHA (c) and CoHA (d).

agreement with the calculated values from DFT calculations which assumed a 10 atomic percent substitution. The small differences in peak position between experimental and theoretical diffraction patterns arise from the differences in the calculated and observed lattice parameters. From atomic absorption spectroscopy, the  $M^{2+}/Ca$  ratios were determined to be 0.07, 0.12, and 0.05 for MnHA, FeHA, and CoHA, respectively. Energy dispersive X-ray spectroscopy was used to assess the stoichiometry of MHA, the  $(Ca + M^{2+})/P$  ratios were determined to be 1.72,

Table 3 Calculated geometry parameters of HA and  $M^{2+}HA$  for site (1) and site (2) substitutions. Experimental lattice parameters were obtained from Rietveld refinements. Units (Å)

System	Lattice constant	Site (1)	Site (2)	Exp
HA	<i>a</i>	9.55	—	9.4220(1)
	<i>c</i>	6.94	—	6.8867(8)
MnHA	<i>a</i>	9.32	9.47	9.4205(1)
	<i>c</i>	6.81	6.9	6.8805(9)
FeHA	<i>a</i>	9.52	9.38	9.4212(3)
	<i>c</i>	6.85	6.95	6.8841(2)
CoHA	<i>a</i>	9.50	9.52	9.4197(2)
	<i>c</i>	6.85	6.94	6.8820(1)

1.71, and 1.69 for MnHA, FeHA, and CoHA, which is in good agreement with the theoretical Ca/P ratio of 1.66, indicating that each  $M^{2+}$  ion replaces a  $Ca^{2+}$  ion. Changes in lattice parameters due to substitution are typically a function of the substituent percentage. Since the transition metal percentage was not controlled, it is difficult to directly compare the effect between transition metals on the lattice parameters.

Chemical analysis indicates that the transition metals successfully replaced calcium in the HA lattice. Between the two distinct calcium sites, DFT calculations indicate that the Ca(II) site is the most energetically favorable site for all calculated substitutions. Experimentally site occupancies cannot be reliably determined since diffraction data was not collected using synchrotron X-rays or with neutron diffraction. However, previous experimental and theoretical studies of Sr, Pb, and Cd substituted HA systems utilizing synchrotron radiation have shown that DFT does accurately predict the most energetically favorable site occupancy, where Ca(II) site is normally replaced by metal ions.<sup>24–26</sup>

To further understand the effect of transition metals on the structure of hydroxyapatite, DFT calculations were performed on the  $Ca_9M(PO_4)_6(OH)_2$  (MHA) systems. Since the valence state of  $Ca^{2+}$  and  $M^{2+}$  are the same, charge compensation defects were not considered for  $M^{2+}$  substitution at the Ca(I) or Ca(II) sites. Table 3 displays the calculated lattice parameters of the MHA, which are in good agreement with data obtained from experiment. The contraction in the lattice parameters of MHA are attributed to the smaller ionic radius and the more electro-negative character of  $Mn^{2+}$ ,  $Fe^{2+}$ , and  $Co^{2+}$  compared to  $Ca^{2+}$ . These differences result in small structural distortions which is seen in the analysis of the interatomic distances of FeHA, shown in Fig. 1. The observed decrease of the bond length between the Fe and the first neighboring oxygen atoms results from iron being more electronegative than calcium attracting the oxygen atoms closer. The Fe at the Ca(2) results in one broke bond with oxygen, where the O(1) bond is broken.

The FTIR spectra of the pure and substituted HA are displayed in Fig. 4 along with the corresponding spectra computed using DFT. It can be seen that both results are in good agreement, indicating that the nearest neighbor environments and bond strengths are correctly captured in the DFT computations. There is a difference between the computed and experimental spectra. The error between the experimental and theoretical spectra for HA, MnHA, FeHA, CoHA, are 3.4%, 3.3%, 3.8%, and 4.1%, respectively. Overall, the theoretical and experimental



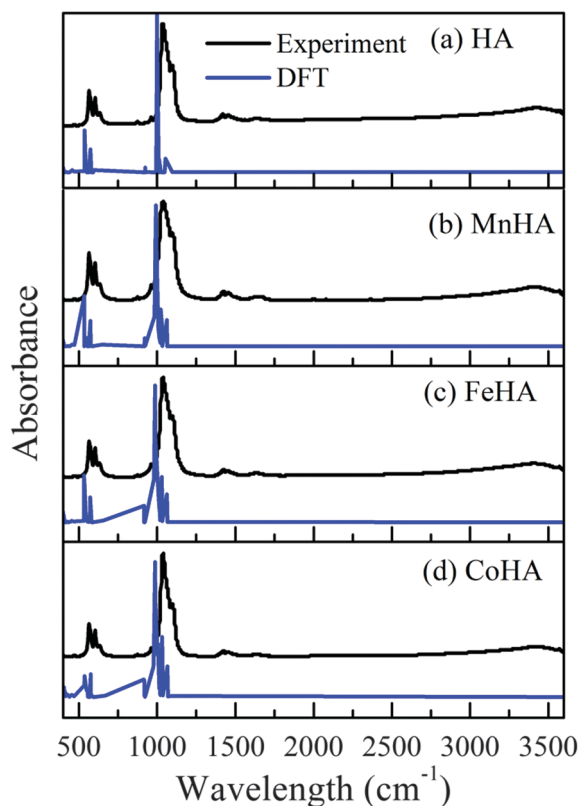


Fig. 4 FTIR spectra of the synthesized HA and MHA. HA (a), MnHA (b), FeHA (c), and CoHA (d).

spectra are in good agreement. Bhat *et al.* observed error ranging from 2% to 8% between theoretical and experimental spectra for stoichiometric HA.<sup>35</sup> Also, the theoretical model of a perfect crystal does not take into account defects such as vacancies that are often present in real samples. Bhat *et al.* calculated the IR spectra of calcium deficient HA and observed a shift in the vibrational modes of phosphate to lower wavenumbers.<sup>35</sup> The MHA spectra display typical vibrational modes of HA, indicating successful substitution of  $M^{2+}$  into the HA lattice. The vibrational modes of phosphate are observed around 1040 and 1095  $\text{cm}^{-1}$  corresponding to the  $\nu_3$  mode, and the  $\nu_4$ ,  $\nu_1$ , and  $\nu_2$  modes are present around 570, 600, and 480  $\text{cm}^{-1}$ , respectively. The vibrational and liberation absorption bands for  $\text{OH}^-$  are observed around 3550 and 640  $\text{cm}^{-1}$ , respectively, which are not predicted in the theoretical spectra. Other absorption bands observed in the experimental spectra that are not predicted by theory are the absorption bands observed at 1450, 1410, and 865  $\text{cm}^{-1}$ , which correspond to the  $\nu_3$  and out of plane  $\nu_2$  mode for  $\text{CO}_3^{2-}$ . The presence of the  $\text{CO}_3^{2-}$  bands is a result of the HA synthesis procedure at ambient atmosphere. Carbonated HA is formed from the adsorption of  $\text{CO}_2$  from the ambient atmosphere. From the position of the  $\text{CO}_3^{2-}$  bands the substitution of  $\text{CO}_3^{2-}$  is predominately at the phosphate site.

High resolution XPS spectra of the substituted transition metal in MHA, as shown in Fig. 5(a, c and e) and tabulated in Table 4, displays two peaks which are attributed to the  $2p_{3/2}$  and  $2p_{1/2}$  states from spin-orbit splitting.<sup>49</sup> For MnHA the peaks at 643.4 and 654.6 eV are the  $2p_{3/2}$  and  $2p_{1/2}$  peak pairs from the

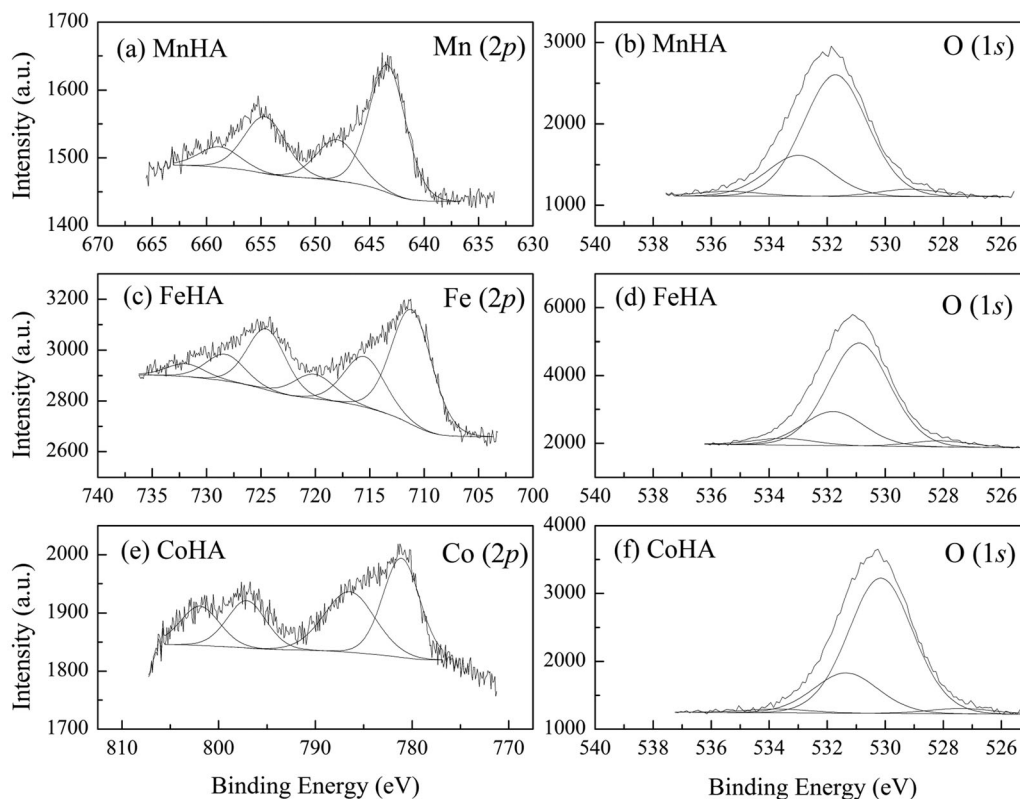


Fig. 5 XPS spectra of M 2p (left) and O 1s (right) peaks. MnHA (a and b), FeHA (c and d) and CoHA (e and f).

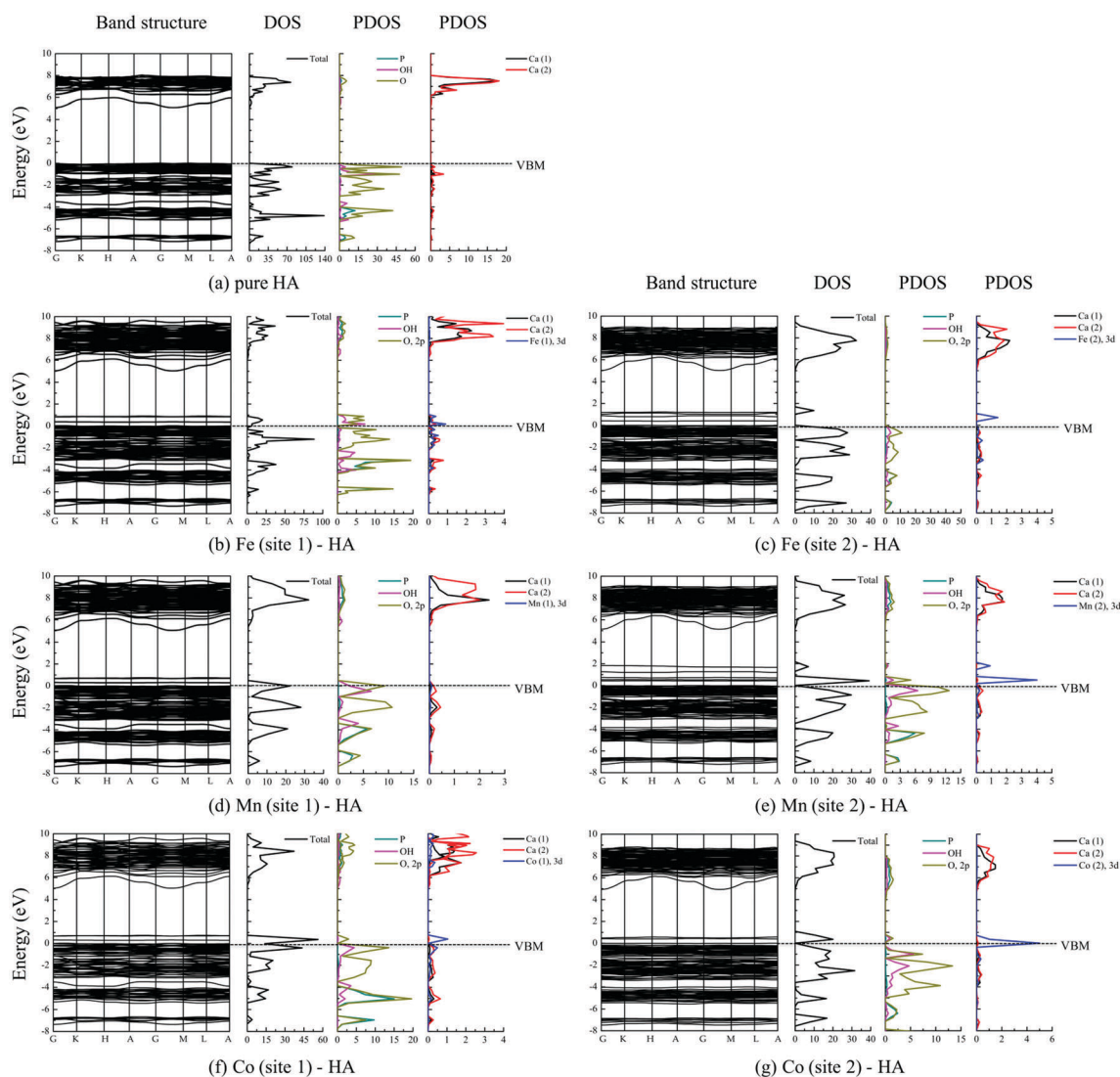
**Table 4** Peak positions from deconvoluted MHA transitions metal (M) 2p spectra in eV

	M	
	2p <sub>3/2</sub>	2p <sub>1/2</sub>
MnHA	643.4, 648.0	654.6, 658.8
FeHA	711.1, 715.4, 720.1	724.5, 728.3, 731.9
CoHA	781.1, 786.5	797.1, 801.9

spin-orbit splitting, which corresponds with the Mn<sup>2+</sup> substituted for Ca in the HA lattice.<sup>50</sup> The peaks at 648.0 and 658.8 eV are also 2p peak pairs, but are shakeup satellites of Mn<sup>2+</sup>. In the case of CoHA the 2p peak pair at 781.1 and 797.1 eV corresponds to Co<sup>2+</sup> substituted in either Ca site in the HA lattice and the 2p peak pair at 786.5 and 801.9 eV are from shakeup satellites. The strong nature of the shakeup satellite is characteristic of Co<sup>2+</sup> with paramagnetic properties.<sup>51</sup> The peaks at 711.1 and 720.1 eV in the FeHA spectra are the 2p peak pairs from spin-orbit splitting

characteristic of Fe<sup>2+</sup>PO<sub>4</sub> bonding, which corresponds to Fe<sup>2+</sup>HA in the lattice.<sup>52</sup> The 2p peak pair around 715.4 and 724.5 eV can be attributed to iron residing on surface sites, which requires more energy to create a photoelectron.<sup>52</sup> The last peak pair around 720.1 and 731.9 eV is a shakeup satellite of iron. The O 1s spectra were also acquired and all MHA displayed peaks around 528, 530, 531, and 533 eV. It should be noted that the peak at 531 eV is characteristic of O in HA (NIST database) and the peak around 530 eV is characteristic of surface carbonate (NIST database). The peak at 533 eV is typical of bridging oxygen's. The peak at 528 eV can be attributed to C–O bonding from the carbon tape used for mounting powders. No recognizable metal oxide bonding was identified in the transition metal 2p or O 1s spectra. All the paramagnetic transition metals appear to have successfully incorporated into the HA lattice in a divalent state.

XPS is successful at probing the electronic structure through core level electrons but does not provide information about the interactions in the valance band. To address this point,

**Fig. 6** Projected density of states for MHA.

**Table 5** Band gap of MHA, given in eV. The data in parenthesis is experimental results from literature

System	Site (1)	Site (2)
HA	5.09	5.3 <sup>a</sup>
MnHA	4.29	3.3
FeHA	4.15	3.8
CoHA	4.29	4.36

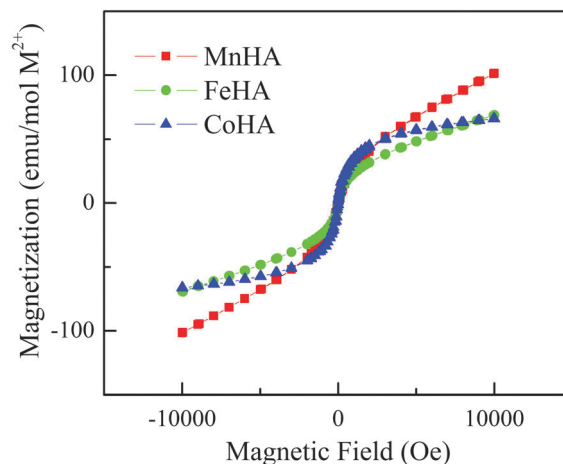
<sup>a</sup> Ref. 53.

we calculated the band structure and density of states, shown in Fig. 6, using DFT to further understand the effect of transition metals on the electronic properties of HA. The valance band maximum of HA is set as 0 eV in all figures, and the calculated PBE band gap of pure HA is 5.09 eV which is consistent with the previously computed result of 5.30 eV from Matsunaga *et al.*<sup>53</sup>

When the transition metals are substituted into HA, the band gap of the MHA decreases compared to pure HA as shown in Fig. 6 and Table 5. In order to describe the influence of specific atoms on the electronic structure of HA, density of states (DOS) and projected density of states analysis were performed, as shown in Fig. 6. Comparing to the partial DOS of pure HA, new energy states of ion-substituted HA are introduced at the top of the valance band owing to the interaction between 2p states of O and the 3d energy states of Fe, Co, and Mn. The interaction between these states results in decreasing the band gap in MHA.

Also, when Fe or Mn is substituted at the Ca(2) site, the band gap is smaller than substitution at the Ca(1) site. The structural distortion from Fe, Mn, and Co at site (1) is less compared to site (2) substitution. The coordination number of Fe, Mn, and Co at site (1) is the same as that of Ca(1), which is 9. Small variations of structure leads to extension of the band tail. However, the interaction between Fe (or Mn) at site (2) and the OH ion lead to the rotation of the OH ion, which results in the production of discrete states in the band structure of HA.

Magnetization measurements depicted in Fig. 7 show that all the substituents induced a transition from diamagnetic behavior, which is characteristic of HA, to a superparamagnetic behavior. A small ferromagnetic signal is observed at low magnetic fields followed by a strong paramagnetic response at high fields. The observed superparamagnetic response may be due to  $M^{2+}$  residing in Ca(2) sites or surface effects as previously discussed.<sup>18</sup> In brief, as the surface to volume ratio increases, a larger number of atoms reside on the surface possibly resulting in dangling bonds, a change in coordination and surface disorder, which can lead to an increase in their magnetization values.<sup>54</sup> The transition from diamagnetic to paramagnetic response is in agreement with the theoretical calculations. For a 10 percent substitution on the Ca-site, our theoretical modelling predicts magnetic moments of 5, 4 and 3 Bohr magnetons for MnHA, FeHA, and CoHA, respectively. Since the MHA fabricated by different approaches have varying levels of  $M^{2+}$  substituted into the lattice, magnetization curves were normalized to the molar amount of  $M^{2+}$  per MHA. In agreement with the calculations, MnHA displayed the highest magnetization value of  $100 \text{ emu mol}^{-1} M^{2+}$  at 10 kOe (or 1 Tesla) of all the substituents, while FeHA and CoHA had similar



**Fig. 7** Mass magnetization of MnHA, FeHA, and CoHA normalized to the transition metal content.

magnetization values around  $65 \text{ emu mol}^{-1} M^{2+}$ . Qualitatively, the slope of the high field magnetization data is also indicative of the magnitude of the magnetic moment and the observed slopes are  $\text{MnHA} > \text{FeHA} > \text{CoHA}$ , which is in good agreement with the theoretical calculations.

## 5. Conclusions

Manganese, iron and cobalt were successfully substituted into hydroxyapatite (HA) and the observed crystal structure, FTIR spectra, electronic and magnetic properties were in good agreement with DFT calculations. The presence of transition metals in a +2 oxidation state was confirmed by XPS, FTIR, and XRD identified the MHA without any detectable impurity phases. Incorporation of magnetic ions into the HA lattice was sufficient to induce a transition from the diamagnetic behavior (for pure HA) to a paramagnetic one, with the observed magnetization values in the order of  $\text{MnHA} > \text{FeHA} > \text{CoHA}$ .

The results from DFT calculations were in good agreement with experimental data. Thus, the DFT calculation could be a useful tool for predicting physical and magnetic properties of different ion-substituted hydroxyapatite powders, and such a tool can guide future studies on the synthesis of ion-substitute hydroxyapatite with desired properties.

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